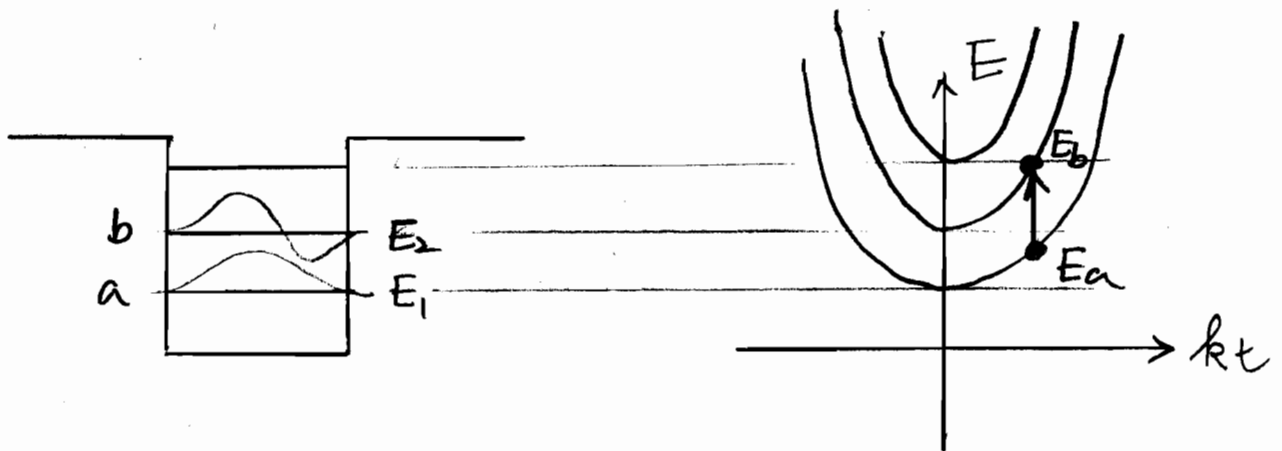


# Intersubband Absorption



Infrared detector (3~5  $\mu\text{m}$ , 8~10  $\mu\text{m}$ )

Gain: Quantum Cascade laser (QCL)

$$\begin{cases} \psi_a(\vec{r}) = u_c(\vec{r}) \cdot \frac{e^{i\vec{k}_c \cdot \vec{r}}}{\sqrt{A}} \phi_1(\delta) \\ \psi_b(\vec{r}) = u_{c'}(\vec{r}) \cdot \frac{e^{i\vec{k}_{c'} \cdot \vec{r}}}{\sqrt{A}} \phi_2(\delta) \end{cases}$$

Note that  $\langle u_c | u_{c'} \rangle \approx 1$ ,  $\langle \phi_2 | \phi_1 \rangle = 0$   
 (In inter-band case  $\langle u_c | u_v \rangle = 0$ )

$$\begin{aligned} \vec{\mu}_{ba} &= \langle \psi_b(\vec{r}) | e\vec{r} | \psi_a(\vec{r}) \rangle \\ &\cong \langle u_c | u_{c'} \rangle \cdot \underbrace{\int \frac{e^{-i\vec{k}_{c'} \cdot \vec{r}}}{\sqrt{A}} \frac{e^{i\vec{k}_c \cdot \vec{r}}}{\sqrt{A}} \phi_2^*(\delta) \cdot e\vec{r} \cdot \phi_1(\delta) d\vec{r}} \end{aligned}$$

$$\approx \delta_{\vec{k}_c, \vec{k}_{c'}} \cdot \langle \phi_2 | e\vec{r} | \phi_1 \rangle$$

$$= \delta_{\vec{k}_c, \vec{k}_{c'}} \langle \phi_2 | e\delta | \phi_1 \rangle \cdot \hat{\delta}$$

$$= \delta_{\vec{k}_c, \vec{k}_{c'}} \cdot \mu_{21} \hat{\delta}$$

anti-symm in  $\delta$   
 $\langle \phi_2 | x | \phi_1 \rangle = 0$   
 $\langle \phi_2 | y | \phi_1 \rangle = 0$

$$E_a = E_1 + \frac{\hbar^2 k_t^2}{2m_e^*}$$

$$E_b = E_2 + \frac{\hbar^2 k_t^2}{2m_e^*}$$

$$g(\Delta E) = \frac{\left(\frac{\Gamma}{2}\right)}{(\Delta E)^2 + \left(\frac{\Gamma}{2}\right)^2} \quad \text{Lorentzian Lineshape}$$

$$\alpha(\hbar\omega) = \frac{\omega}{n_r c \epsilon_0} \cdot \frac{2}{V} \sum_{\vec{k}_t} \sum_{\vec{k}_t'} g(E_b - E_a - \hbar\omega) \cdot \underbrace{|\hat{e} \cdot \vec{\mu}_{ba}|^2}_{\delta_{\vec{k}_t, \vec{k}_t'} \cdot |\mu_{21}|^2} \cdot (f_a - f_b)$$

$$\hat{e} \cdot \vec{\mu}_{ba} = \hat{e} \cdot \hat{z} \cdot \mu_{21} \cdot \delta_{\vec{k}_t, \vec{k}_t'}$$

None zero when  $\hat{e} = \hat{z}$  (TM)

$$\alpha(\hbar\omega) = \frac{\omega}{n_r c \epsilon_0} \cdot \frac{2}{V} \sum_{\vec{k}_t} g(E_2 - E_1 - \hbar\omega) \cdot |\mu_{21}|^2 \cdot (f_a - f_b)$$

↓  
indep of  $\vec{k}_t$  ←

$$= \frac{\omega}{n_r c \epsilon_0} \cdot g(E_2 - \hbar\omega) \cdot |\mu_{21}|^2 \cdot \underbrace{\frac{2}{V} \sum_{\vec{k}_t} (f_a - f_b)}_{N_1 - N_2}$$

$$= \frac{\omega}{n_r c \epsilon_0} \cdot g(E_2 - \hbar\omega) \cdot |\mu_{21}|^2 \cdot (N_1 - N_2)$$

$$N_i = \frac{m_e^* k_B T}{\pi \hbar^2 L_z} \cdot \ln \left( 1 + e^{\frac{E_F - E_i}{k_B T}} \right) \quad i=1, 2$$

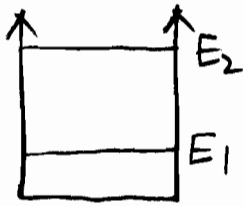
Low temp limit:  $(E_F - E_i) \gg k_B T$

$$N_i \approx \frac{m_e^*}{\pi \hbar^2 L_z} \cdot (E_F - E_i) \quad i=1, 2$$

When  $E_F \ll E_2$ ,  $N_2 \ll N_1$ ,  $\alpha(\hbar\omega) \propto N_1$  63

When  $E_F \gg E_2$ ,  $\alpha(\hbar\omega) \propto (N_1 - N_2) \propto (E_2 - E_1) = \text{const.}$

$$\alpha(\hbar\omega) = \frac{\omega}{n_c c \epsilon_0} g(E_2 - \hbar\omega) |u_{21}|^2 \left( \frac{m_e^*}{\pi \hbar^2 L_3} \right) (E_2 - E_1)$$



0 10 nm

GaAs

$$m_e^* = 0.067 m_0$$

$$E_1 = \frac{\hbar^2}{2m_e^*} \left( \frac{\pi}{L_3} \right)^2 = 56.5 \text{ meV}$$

$$E_2 = 4E_1 = 226 \text{ meV}$$

$$\phi_1 = \sqrt{\frac{2}{L_3}} \cdot \sin\left(\frac{\pi}{L_3} z\right)$$

$$\phi_2 = \sqrt{\frac{2}{L_3}} \sin\left(2\frac{\pi}{L_3} z\right)$$

$$u_{21} = e \int_0^{L_3} \phi_2(z) \cdot z \cdot \phi_1(z) dz$$

$$= \frac{2e}{L_3} \int_0^{L_3} \sin\left(\frac{\pi}{L_3} z\right) \cdot z \cdot \sin\left(2\frac{\pi}{L_3} z\right) dz$$

$$= -\frac{16}{9\pi^2} e L_3$$

If  $N = 10^{18} \text{ cm}^{-3}$

Assume  $E_1$  is occupied.  $E_2$  not.

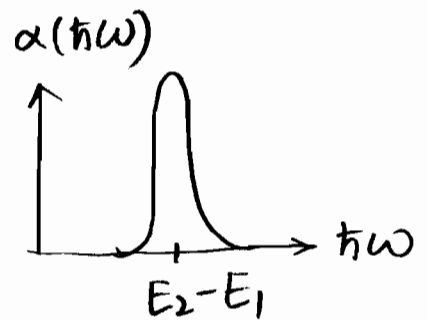
$$E_F - E_1 = N \frac{\pi \hbar^2 L_3}{m_e^*} = 36 \text{ meV}$$

Indeed  $E_2$  is not occupied.

Peak absorption

$$\hbar\omega = E_2 - E_1 \cong 170 \text{ meV}$$

$$\lambda = \frac{1.24}{0.17} = 7.3 \mu\text{m}$$



$$\alpha_{\text{peak}} = \frac{\omega}{n_r c \epsilon_0} \cdot \frac{|M_{21}|^2}{\left(\frac{\Gamma}{2}\right)} \underbrace{(N_1 - N_2)}_{\substack{\text{ss} \\ N}} \cong 10^4 \text{ cm}^{-1}$$

Note = absorption  $\propto N$