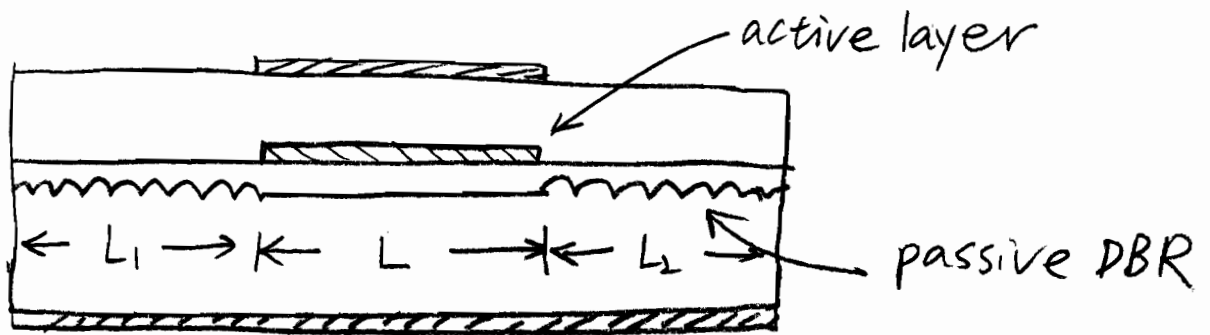


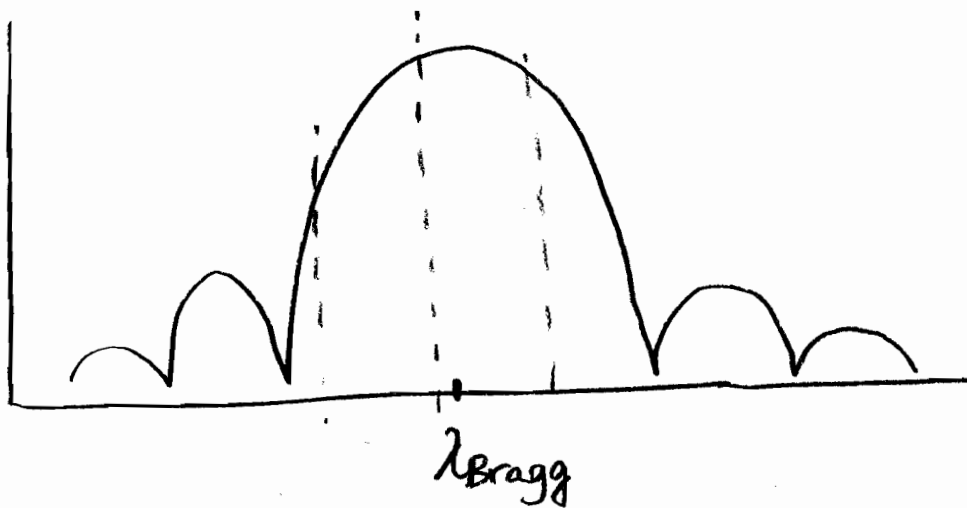
# Distributed Bragg Reflector (DBR) Laser

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\* Optical cavity formed by passive gratings (DBR's)



\* Wavelength-dependent reflectivity



- Lasing wavelength  $\approx \lambda_{\text{Bragg}} = 2 \cdot n_{\text{eff}} \cdot \Lambda$
- Exact wavelength is determined by modes  
↳ With 0.5 mode spacing from the peak reflectivity

Threshold condition:

$$r_1 r_2 e^{2\beta \cdot L} = 1 \quad \text{round-trip condition for field.}$$

$$\beta = \beta_0 - i \frac{1}{2} (\Gamma g - \alpha)$$

$$r_1 = \sqrt{R_1} e^{i\phi_1}$$

$$r_2 = \sqrt{R_2} e^{i\phi_2}$$

$$\Rightarrow \beta_{th} = \frac{1}{l} (\alpha + \alpha_m) = \frac{1}{l} (\alpha + \frac{1}{2L} \ln \frac{1}{R_1 R_2})$$

$$2\beta_0 L + \phi_1 + \phi_2 = 2m\pi$$

$$r_1 = |r_1| e^{i\phi_1}$$

$$= \frac{-k^* \sinh SL_1}{\Delta \beta \sinh SL_1 + i S \cosh SL_1}$$

$$\beta = iS = \sqrt{\delta^2 - |k|^2} = i \sqrt{|k|^2 - \delta^2}$$

$$S = \sqrt{|k|^2 - \delta^2}$$

Note  $\phi_1$  is a function of  $\beta_0$  also.

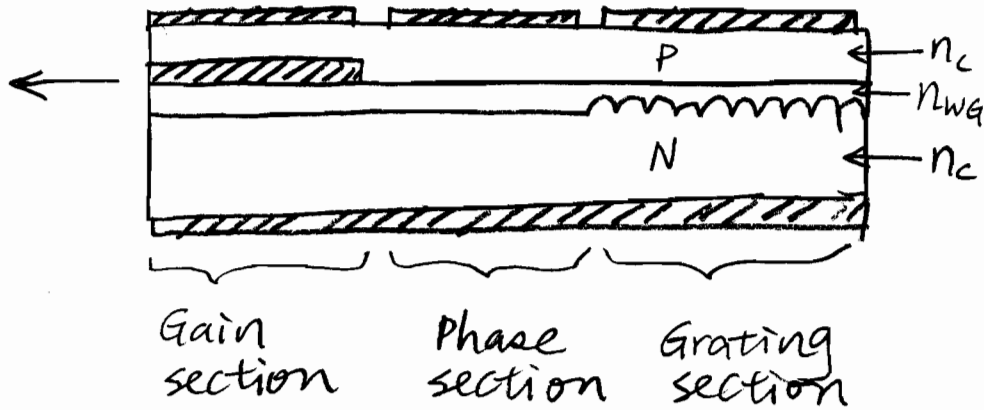
$$\frac{d\phi_1}{d\beta} = 2 \cdot L_{eff,1}$$

Likewise  $\frac{d\phi_2}{d\beta} = 2 \cdot L_{eff,2}$

$$\Rightarrow 2\beta_0 (L + L_{eff,1} + L_{eff,2}) = 2m\pi$$

\* Mode spacing is determined by  
active section length + effective grating lengths

\* A unique benefits of DBR  $\rightarrow$  Tunable laser



$$\lambda_{\text{Bragg}} = 2 \cdot n_{\text{eff}} \cdot \Lambda$$

↑  
effective index of grating section.

$$n_{\text{eff}} = \Gamma n_{\text{wg}} + (1 - \Gamma) n_c$$

$n_{\text{wg}}$  can be changed by injecting carriers  
(called free-carrier plasma effect)

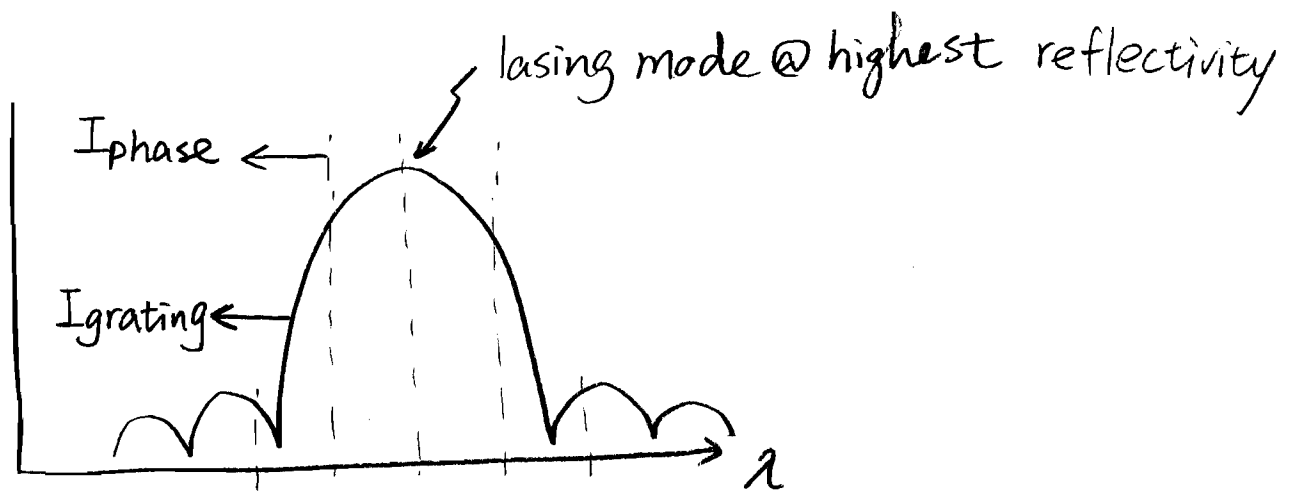
Typically

$$\frac{dn}{dN} = -10^{-20} \text{ cm}^3$$

e.g.  $\Delta N = 10^{18} \text{ cm}^3$

$$n \rightarrow n - 0.01$$

$$\frac{\Delta n}{n} \sim 0.1\% \text{ to } 1\%$$

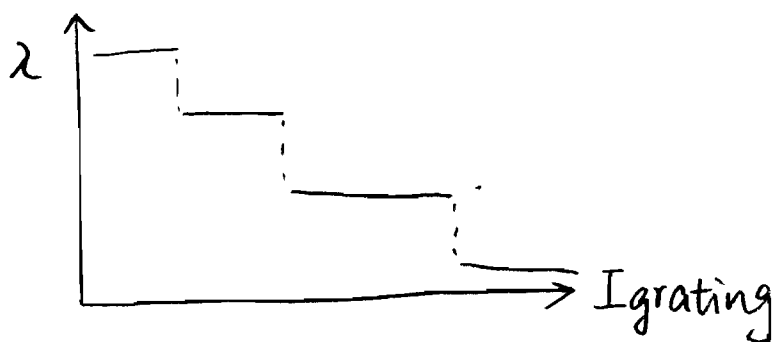


Increasing grating section current ( $I_{\text{grating}}$ ) moves reflectivity curve to shorter  $\lambda$  side, but does not change mode position much

$\Rightarrow$  discrete tuning (step  $\cong$  mode spacing)

Increasing phase section current ( $I_{\text{phase}}$ ) moves modes to shorter  $\lambda$  side, but does not change the reflectivity curve

$\Rightarrow$  continuous tuning with  $\Delta\lambda_{\text{max}} = \text{mode sp.}$



$\Rightarrow$  Combining  $I_{\text{grating}}$  and  $I_{\text{phase}}$

$\Rightarrow$  Continuous tuning over a large range