

Cromatic Lecture # 9

Quantum Teleportation Using Photons

Based on Bouwmeester, et al., Nature 390, 575 (1997) [zeilinger97].

Idea: Alice has qubit state $|4\rangle = \alpha|0\rangle + \beta|1\rangle$ which she wants to send to Bob with the help of 2 other entangled qubits and some classical information.



Qubit System: Polarization state of photons,

Basis states: \vec{E} -Horizontal $= |0\rangle$, \vec{E} -Vertical $= |1\rangle$
 \vec{E} -field polarization states.

\Rightarrow Photon having $|4\rangle = \alpha|0\rangle + \beta|1\rangle$ has probability $|\alpha|^2$ to be found in Horiz. polarization state and $|\beta|^2$ prob. to be found in vertical polariz. state

\Rightarrow Need 3 photons (3 qubits)

Photon #1: $|4\rangle_1 = \alpha|0\rangle_1 + \beta|1\rangle_1$, belongs to Alice.
This is the state she wants to teleport.

Photons #2 & #3: $|4\rangle_{23} = |\Psi_{in}\rangle = \frac{1}{\sqrt{2}}(|0\rangle_2|1\rangle_3 - |1\rangle_2|0\rangle_3)$

This is an entangled pair. Alice has one of the photons (Photon #2) and Bob has the other (Photon #3).

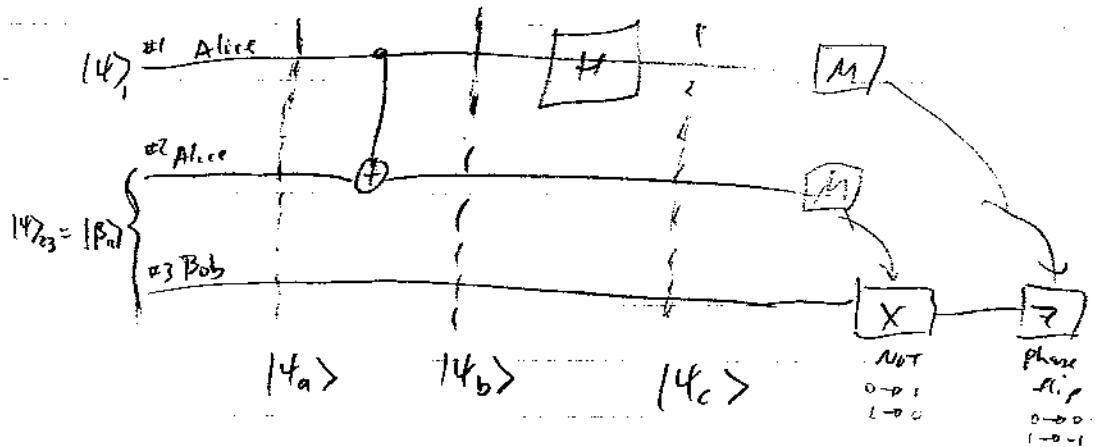
\Rightarrow Alice wants to teleport the state of photon #1 to photon #3 in Bob's possession.

→ How it works in theory

$$\text{Inputs} \rightarrow \begin{cases} |\psi_1\rangle = \alpha|10\rangle + \beta|11\rangle \\ |\psi_2\rangle = |\beta_{11}\rangle = \frac{1}{\sqrt{2}}(|101\rangle - |110\rangle) \end{cases}$$

Consider Teleportation
Circuit we saw
before in class

(2)



$$|\psi_a\rangle = |\psi_1\rangle \otimes |\beta_{11}\rangle = (\alpha|10\rangle + \beta|11\rangle) \cdot \frac{1}{\sqrt{2}}(|101\rangle - |110\rangle)$$

$$\Rightarrow |\psi_a\rangle = \frac{1}{\sqrt{2}}\alpha|10\rangle(|101\rangle - |110\rangle) + \frac{1}{\sqrt{2}}\beta|11\rangle(|101\rangle - |110\rangle)$$

$$\Rightarrow |\psi_b\rangle = \frac{1}{\sqrt{2}}\alpha|10\rangle(|101\rangle - |110\rangle) + \frac{1}{\sqrt{2}}\beta|11\rangle(|111\rangle - |100\rangle)$$

$$\Rightarrow |\psi_c\rangle = \frac{1}{2}\alpha(|10\rangle + |11\rangle)(|101\rangle - |110\rangle) + \frac{1}{2}\beta(|10\rangle - |11\rangle)(|111\rangle - |100\rangle)$$

$$= \frac{1}{2}[\alpha|1001\rangle - \alpha|1010\rangle + \alpha|1101\rangle - \alpha|1110\rangle + \beta|1011\rangle - \beta|1100\rangle - \beta|1111\rangle + \beta|1000\rangle]$$

$$= \frac{1}{2}[|100\rangle[|\alpha|1\rangle - \beta|0\rangle] - |101\rangle[|\alpha|10\rangle - \beta|11\rangle] + |110\rangle[|\alpha|11\rangle + \beta|10\rangle] - |111\rangle[|\alpha|10\rangle + \beta|11\rangle]]$$

→ Useful thing?

Now Do measurement of qubits #1 & #2 $\xrightarrow{\text{in standard basis}}$ Act on Qubit #3
depending on results of this measurement

(3)

→ Suppose measurement of qubits #1 & #2 registers "1" and "1"
 (i.e., $|11\rangle$)
 $\Rightarrow |14\rangle_3 = \alpha|10\rangle + \beta|11\rangle$ since $|14\rangle \propto |11\rangle\langle 11|4\rangle$

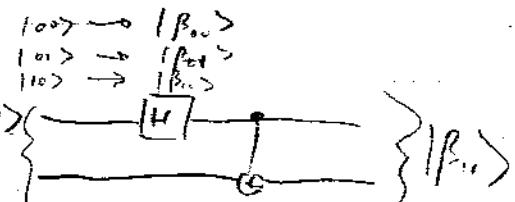
In this case the qubit is teleported and Bob doesn't have to do anything to qubit #3... (no "X" or "Z")

This is the situation that has been performed experimentally with photons. [other cases haven't been done yet]
 to best of my knowledge

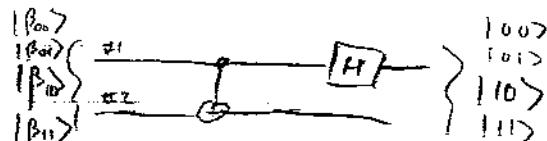
But it's done w/ a small conceptual trick

Notice:

This gate transforms the standard basis into ^{the} Bell state basis



And so reversed gate transforms Bell states into standard basis



Note: This is exactly what we have in the teleportation circuit. This means that if qubits #1, #2 are projected onto $|B_{11}\rangle$ before hitting this gate and the gate is removed, then $|B_{11}\rangle$ will play exactly same role that $|11\rangle$ would have played if gate stayed. We still get teleportation to qubit #3, only now it's tensored w/ $|B_{11}\rangle$ instead of $|11\rangle$. Basically we just replace role of $|11\rangle$ w/ $|B_{11}\rangle$.

So, instead of running qubits #1 & #2 through ~~f~~ and projecting them onto $|11\rangle$ at the output, why not just directly project qubits #1 & #2 onto $|B_{11}\rangle$ directly at the beginning?

[easy to show it's same mathematically, see supplementary std] (4)

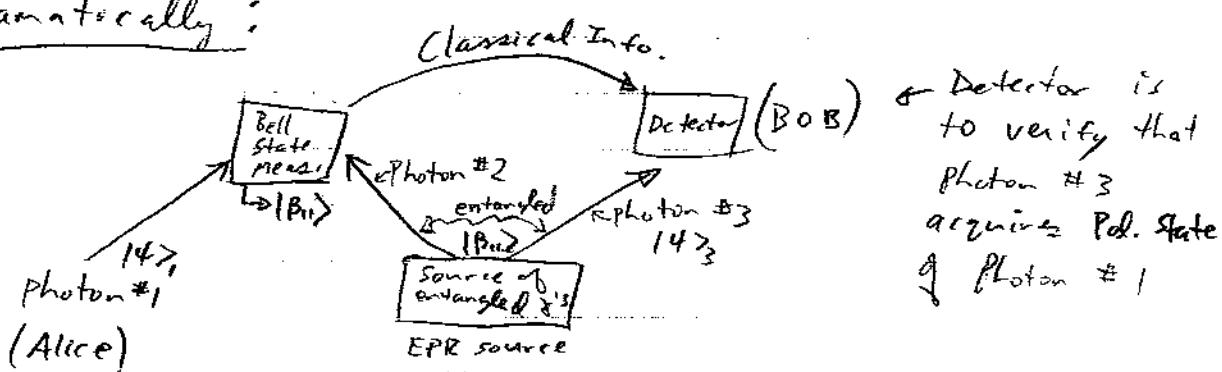
This is what is done experimentally with photons to achieve Quantum teleportation. A "Bell state measurement" is performed which projects qubits #1 & #2 onto $|B_{11}\rangle$. Alice's state is then ~~teleported~~ onto qubit #3 without any further work.

So, the new Q.Teloprtation scheme is as follows : Alice performs a Bell state measurement on photons #1 & #2 and projects them onto the Bell state $|B_{11}\rangle = |1\rangle_1|0\rangle_2 - |1\rangle_1|1\rangle_2$

This forces photon #3 to be in the former state of photon #1 ($|1\rangle_3$)

$$\Rightarrow |1\rangle_3 \rightarrow \alpha|0\rangle + \beta|1\rangle$$

Diagrammatically :



Classical info. is needed to verify that photons #1 and #2 have been projected onto the correct Bell state ($|B_{11}\rangle$). The teleportation will only work for projection on to this state [i.e. this is the one choice out of 4 where Bob doesn't have to change the state of Photon #3 to obtain $|1\rangle_3$]

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So, how do you do this in real life?

There are 2 important experimental tricks:

1) ^{Must} Create 2 entangled photons in $|\beta_{11}\rangle$ state:

$$|\beta_{11}\rangle = \frac{1}{\sqrt{2}} (|H\rangle_{\text{photon}1}^{\#1} |V\rangle_{\text{photon}2}^{\#2} - |V\rangle_{\text{photon}1}^{\#1} |H\rangle_{\text{photon}2}^{\#2})$$

$$|H\rangle \rightarrow |0\rangle, \quad |V\rangle \rightarrow |1\rangle$$

$$\Rightarrow |\beta_{11}\rangle = \frac{1}{\sqrt{2}} (|01\rangle - |10\rangle)$$

2) ^{Must} Perform Bell state measurement on 2 photons. Project 2 photons onto $|\beta_{11}\rangle$ state. How?

$|^{st}$ let's talk about how to
create entangled photons;

(6)

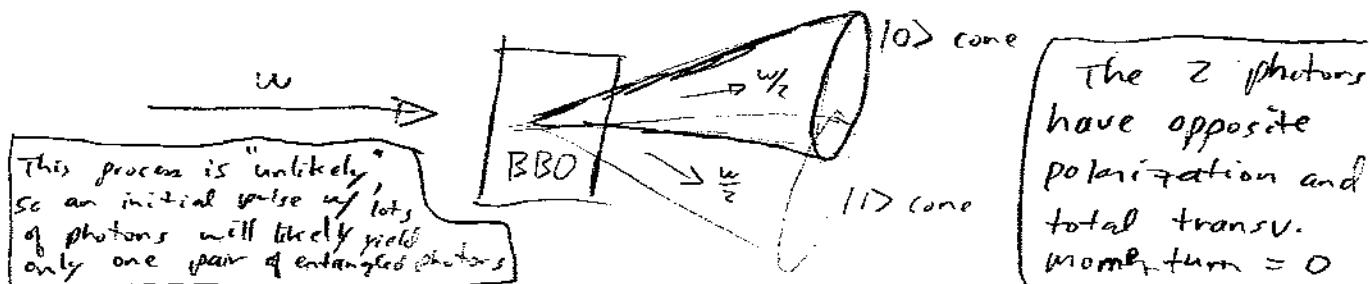
① How do you get an EPR pair of photons? i.e., an entangled photon state?

Answer: Use Nonlinear crystal capable of "parametric down conversion"

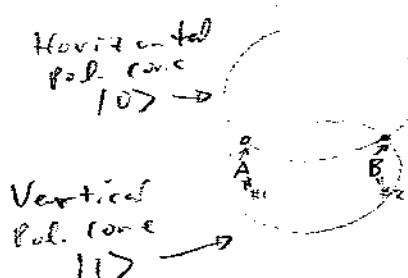
Special Nonlinear X'tals (such as barium borate) have the property that you can send in 1 photon at freq. ω and get out 2 photons at freq. $\omega/2$.

The 2 photons come out in 2 cones.

One cone has a Horiz. polarized photon $|10\rangle$ and the other has a vert. polarized photon $|11\rangle$



Entangled photons come from capturing photons at the intersection of the cones:



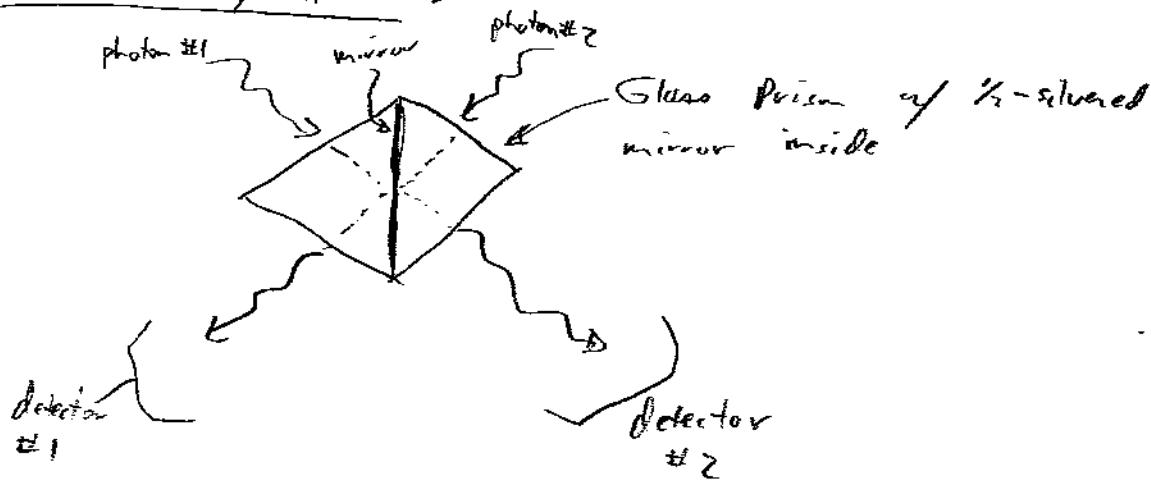
The 2 photons coming from points A and B have indeterminate polarization but they must have opposite polarization \Rightarrow they are entangled:

$$|4\rangle_{AB} = \frac{1}{\sqrt{2}} (|10\rangle_A |11\rangle_B - |11\rangle_A |10\rangle_B)$$

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(7) How do you perform Bell State measurement?

Use Beam splitter!



Assuming photons^{#1 & #2} are incident on the beam splitter at exactly the same time \Rightarrow there are 4 possible polarization states in the

BELL STATE BASIS : $|\beta_{00}\rangle = \frac{1}{\sqrt{2}}(|H\rangle|H\rangle + |V\rangle|V\rangle) = \frac{1}{\sqrt{2}}(|100\rangle + |111\rangle)$

[initial 2 photon state
is a superposition
of these \downarrow]

$$|\beta_{01}\rangle = \frac{1}{\sqrt{2}}(|H\rangle|V\rangle + |V\rangle|H\rangle) = \frac{1}{\sqrt{2}}(|101\rangle + |110\rangle)$$

$$|\beta_{10}\rangle = \frac{1}{\sqrt{2}}(|H\rangle|H\rangle - |V\rangle|V\rangle) = \frac{1}{\sqrt{2}}(|100\rangle - |111\rangle)$$

$$|\beta_{11}\rangle = \frac{1}{\sqrt{2}}(|H\rangle|V\rangle - |V\rangle|H\rangle) = \frac{1}{\sqrt{2}}(|101\rangle - |110\rangle)$$

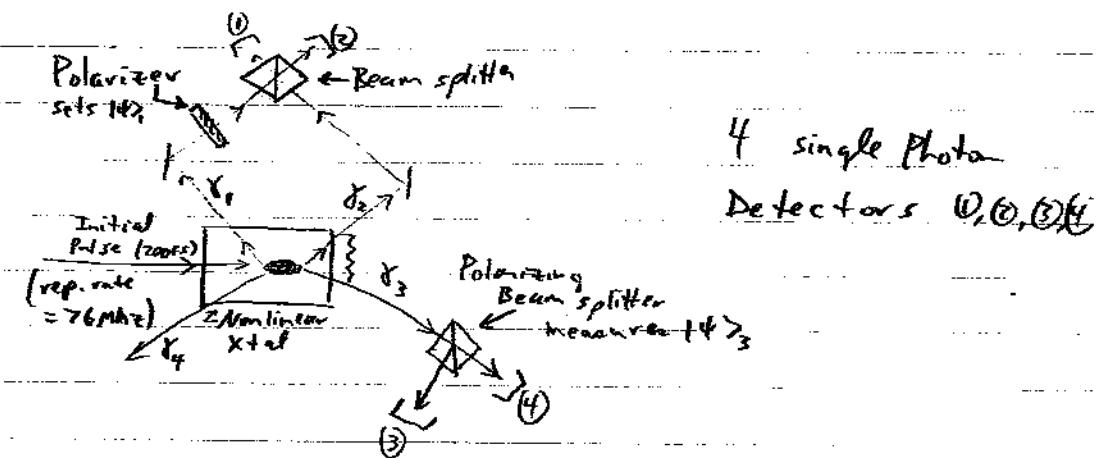
It turns out that the properties of a beam splitter are such that if detectors #1 & #2 click simultaneously \Rightarrow photons are in $|\beta_{11}\rangle$ state. Only one detector or the other would click for the other Bell states.

Showing why this is true takes some tricky Quantum optics (ref. ?) Not discussed here.

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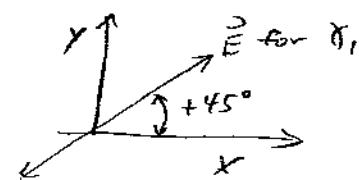
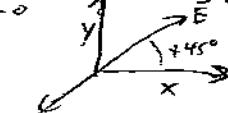
So, how do we actually do the exp. in the lab??

Use laser, parametric down conversion, beam splitters, single photon detection:



(9)

Experiment & Results :

- 1) Initial pulse hits Nonlinear Xtal and creates 2 entangled pairs: $|\gamma_1 \gamma_4\rangle$ and $|\gamma_2 \gamma_3\rangle$. γ_4 is thrown away, and γ_1 is used as qubit #1. $\gamma_1, \gamma_2, \gamma_3, \gamma_4$ created sometime
 - 2) γ_1 is passed through a polarizer that sets $|4\rangle$, at a polarization of $+45^\circ$. This is the state we want to teleport to γ_3 .
 
 - 3) γ_2 and γ_3 are entangled into $|\beta_{11}\rangle$ state = $|4\rangle_{23}$
 - 4) γ_1 and γ_2 hit beam splitter at same time. Bell state measurement is performed by looking for simultaneous firing of detectors #1 and #2. If they fire at the same time then we know that Bell state measurement has been performed and $|4\rangle_{12} = |\beta_{11}\rangle = \frac{1}{\sqrt{2}}(|0\rangle_1|1\rangle_2 - |1\rangle_1|0\rangle_2)$ has been detected (these detectors provide the 2 classical bits of information).
 - 5) The state of γ_3 is measured with polarizing beam splitter. If γ_3 is polarized at $+45^\circ$ then detector (4) fires.
 

 If γ_3 is polarized at -45° then detector (3) fires.
 
- This is a complete basis, so need to measure a predominance of detector #4 clicks to know that state is $\overset{\text{selected}}{A}$,

A successful Teleportation event occurs when detectors (1), (2) and (4) go off at the same time. This shows that qubit states for δ_1 and δ_2 have been projected onto the Bell state $|4\rangle_{12}$ and the qubit state for δ_1 has been transported to δ_3 .

It is also important that detector (3) NOT go off for repeated measurements, so that we know that the state of δ_3 is really an eigenstate of $+45^\circ$, and not just being projected onto it.

Stop here!

What is time

Uncertainty in

order ??

[just have ~~to~~ δ_1, δ_2

measured before δ_3]

Supplementary :

(11)

Teleportation Algebra for Zeilinger Experiment:

$$|4\rangle_1 = \alpha|0\rangle_1 + \beta|1\rangle_1 \quad (\text{State we want to transport to } |4\rangle_3) \quad \text{Bob via}$$

$$|4\rangle_{23} = \frac{1}{\sqrt{2}}(|0\rangle_2|1\rangle_3 - |1\rangle_2|0\rangle_3) \quad (\text{entangled qubits split between Alice and Bob})$$

Initial State of 3 photons:

$$\begin{aligned} |4\rangle_{123} &= (\alpha|0\rangle_1 + \beta|1\rangle_1) \cdot \frac{1}{\sqrt{2}}(|0\rangle_2|1\rangle_3 - |1\rangle_2|0\rangle_3) \\ &= \frac{1}{\sqrt{2}} \left[\alpha|0\rangle_1|0\rangle_2|1\rangle_3 - \alpha|0\rangle_1|1\rangle_2|0\rangle_3 + \beta|1\rangle_1|0\rangle_2|1\rangle_3 \right. \\ &\quad \left. - \beta|1\rangle_1|1\rangle_2|0\rangle_3 \right] \end{aligned}$$

Now, do Bell state Measurement of photons #1 & #2 by projecting them onto $|4\rangle_{12} = \frac{1}{\sqrt{2}}(|0\rangle_1|1\rangle_2 - |1\rangle_1|0\rangle_2)$

$$\Rightarrow \text{Define } |\phi\rangle = |4\rangle_{12} \langle 4| |4\rangle_{123} \quad [\text{Proj. of } |4\rangle_{123} \text{ onto } |4\rangle_{12}]$$

$$\Rightarrow \text{After measurement} \Rightarrow \text{New } |4'\rangle_{123} = \frac{|\phi\rangle}{\sqrt{\langle \phi | \phi \rangle}}$$

So, what is $|\phi\rangle = ?$

$$\begin{aligned} |\phi\rangle &= |4\rangle_{12} \left[\frac{1}{\sqrt{2}}(\langle 0|_1\langle 1|_2 - \langle 1|_1\langle 0|_2) \cdot \frac{1}{\sqrt{2}}(\alpha|0\rangle_1|0\rangle_2|1\rangle_3 - \alpha|0\rangle_1|1\rangle_2|0\rangle_3 + \beta|1\rangle_1|0\rangle_2|1\rangle_3 \right. \\ &\quad \left. - \beta|1\rangle_1|1\rangle_2|0\rangle_3 \right] \end{aligned}$$

$$= |4\rangle_{12} \cdot \frac{1}{2} \cdot (\alpha \cdot 0 - \alpha|0\rangle_3 + \beta \cdot 0 - \beta|0\rangle_3 - \alpha \cdot 0 + \alpha|1\rangle_3 + \beta \cdot 0 - \beta|1\rangle_3)$$

$$|\phi\rangle = \frac{1}{2}|4\rangle_{12} (\alpha|0\rangle_3 + \beta|1\rangle_3)$$

(12)

$$\Rightarrow \langle \phi | \phi \rangle = \frac{1}{4} \Rightarrow \sqrt{\langle \phi | \phi \rangle} = \frac{1}{2}$$

$$\text{So, New } |\psi'\rangle_{123} = \frac{|\phi\rangle}{\sqrt{\langle \phi | \phi \rangle}} = -|\phi\rangle_{12} (\alpha|0\rangle_3 + \beta|1\rangle_3)$$

$$\Rightarrow \text{By inspection we see } |\psi'\rangle_{123} = |\psi'\rangle_{12} |\psi'\rangle_3$$

AND $|\psi'\rangle_3 = \alpha|0\rangle_3 + \beta|1\rangle_3$ Exactly what we wanted!

Alice's original qubit #1 state has teleported to qubit #3.