

# CS-184: Computer Graphics

## Lecture #22: Rigid Body Dynamics

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### Announcement

- Final Project Poster Session
  - Wednesday May 1th, 11:00am-2:00 pm
  - Poster stands and tables provided
  - Laptop videos or demos are highly recommended
  - Limited AC outlets
- Final project reports
  - **Hardcopy** due to me by end of Final Exam.
  - *No time for late submission!*
- Final exam
  - Tuesday, May 14th, 3:00-6:00pm
  - 105 Stanley Hall

# Today

- Rigid-body dynamics
- Articulated systems

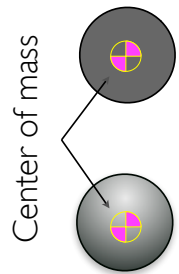
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## A Rigid Body

- **A solid object that does not deform**
- Consists of infinite number of infinitesimal mass points...
- ...that share a single RB transformation
  - Rotation + Translation (no shear or scale)
$$x^W = R \cdot x^L + t$$
  - Rotation and translation vary over time
- Limit of deformable object as  $k_S \rightarrow \infty$

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# A Rigid Body



In 2D:  
Translation 2 “directions”  
Rotation 1 “direction”  
3 DOF Total

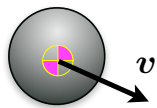
In 3D:  
Translation 3 “directions”  
Rotation 3 “direction”  
6 DOF Total

Translation and rotation are **decoupled**

2D is boring... we'll stick to 3D from now on...

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# Translational Motion



Just like a point mass:

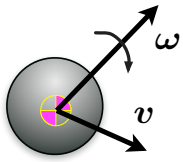
$$\dot{\mathbf{p}} = \mathbf{v}$$

$$\dot{\mathbf{v}} = \mathbf{a} = \mathbf{f}/m$$

Note: Recall discussion on integration...

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## Rotational Motion



Rotation gets a bit odd, as well see...

Rotational “position”  $R$

Rotation matrix

Exponential map

Quaternions

Rotational velocity  $\omega$

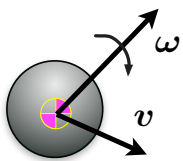
Stored as a vector

(Also called angular velocity...)

Measured in radians / second

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## Rotational Motion



Kinetic energy due to rotation:

“Sum energy (from rotation) over all points in the object”

$$E = \int_{\Omega} \frac{1}{2} \rho \dot{\mathbf{x}} \cdot \dot{\mathbf{x}} du$$

$$E = \int_{\Omega} \frac{1}{2} \rho ([\omega \times] \mathbf{x}) \cdot ([\omega \times] \mathbf{x}) du$$

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# Rotational Motion

$$\begin{aligned}
 H &= \frac{\partial E}{\partial \omega} && H \text{ momentum (angular)} \\
 & && \text{(work)} \\
 H_p &= \frac{\partial E}{\partial \omega_p} \\
 &= \int_{\Omega} \rho \frac{1}{2} (\epsilon_{ijk} \delta_{jp} x_k \epsilon_{iab} \omega_a x_b + \epsilon_{ijk} \omega_j x_k \epsilon_{iab} \delta_{pa} x_b) du \\
 &= \int_{\Omega} \rho \frac{1}{2} (\epsilon_{ipn} x_n \epsilon_{iab} \omega_a x_b + \epsilon_{ijk} \omega_j x_k \epsilon_{iab} \delta_{pa} x_b) du \\
 &= \int_{\Omega} \rho \epsilon_{ipn} x_n \epsilon_{iab} \omega_a x_b du \\
 &= \int_{\Omega} \rho \mathbf{x} \times (\boldsymbol{\omega} \times \mathbf{x}) du \\
 &= \int_{\Omega} \rho \mathbf{x} \times \mathbf{v} du \\
 &= \omega_a \int_{\Omega} \rho \epsilon_{ipn} x_n \epsilon_{iab} x_b du \\
 &= \omega_a \int_{\Omega} \rho (\delta_{pa} \delta_{nb} - \delta_{pb} \delta_{na}) x_n x_b du \\
 &= \omega_a \int_{\Omega} \rho (\delta_{pa} x_n x_n - x_n x_p) du \\
 &= \int_{\Omega} \rho \mathbf{I}_{ap} \omega_a du
 \end{aligned}$$

Angular momentum  
 Similar to linear momentum  
 Can be derived from rotational energy

$$\mathbf{H} = \int_{\Omega} \rho \mathbf{x} \times \dot{\mathbf{x}} du$$

$$\mathbf{H} = \int_{\Omega} \rho \mathbf{x} \times (\boldsymbol{\omega} \times \mathbf{x}) du$$

$$\mathbf{H} = \left( \int_{\Omega} \dots du \right) \boldsymbol{\omega}$$

“Inertia Tensor” not identity matrix...

$$\mathbf{H} = \mathbf{I}\boldsymbol{\omega}$$

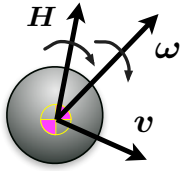
# Inertia Tensor

$$\mathbf{I} = \int_{\Omega} \rho \begin{bmatrix} y^2 + z^2 & -xy & -xz \\ -xy & z^2 + x^2 & -yz \\ -xz & -yz & x^2 + y^2 \end{bmatrix} du$$

See example for simple shapes at <http://scienceworld.wolfram.com/physics/MomentofInertia.html>

Can also be computed from polygon models by transforming volume integral to a surface one.  
 See paper/code by Brian Mirtich.

# Rotational Motion



Conservation of momentum:

$$\mathbf{H}^W = \mathbf{I}^W \boldsymbol{\omega}^W$$

$$\mathbf{H}^W = \mathbf{R} \mathbf{I}^L \mathbf{R}^T \boldsymbol{\omega}^W$$

$$\dot{\mathbf{H}}^W = \dot{\mathbf{R}} \mathbf{I}^L \mathbf{R}^T \boldsymbol{\omega}^W + \mathbf{R} \mathbf{I}^L \dot{\mathbf{R}}^T \boldsymbol{\omega}^W + \mathbf{R} \mathbf{I}^L \mathbf{R}^T \boldsymbol{\alpha}^W$$

Figure is a lie if this really is a sphere...

$$\dot{\mathbf{H}}^W = 0$$

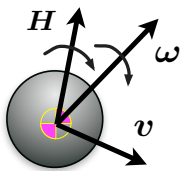
$$\dot{\mathbf{R}} = \boldsymbol{\omega} \times \mathbf{R}$$

$$\boldsymbol{\alpha}^W = (\mathbf{R} \mathbf{I}^L \mathbf{R}^T)^{-1} (-\boldsymbol{\omega}^W \times \mathbf{H}^W)$$

In other words, things wobble when they rotate.

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# Rotational Motion



$$\dot{\mathbf{R}} = [\boldsymbol{\omega} \times] \mathbf{R}$$

$$\dot{\boldsymbol{\omega}} = \boldsymbol{\alpha}$$

Figure is a lie if this really is a sphere...

$$\boldsymbol{\alpha}^W = (\mathbf{R} \mathbf{I}^L \mathbf{R}^T)^{-1} ((-\boldsymbol{\omega}^W \times \mathbf{H}^W) + \boldsymbol{\tau})$$

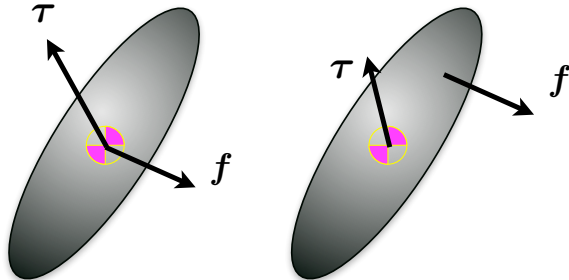
$$\boldsymbol{\tau} = \mathbf{f} \times \mathbf{x}$$

Take care when integrating rotations, they need to stay rotations.

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# Couples

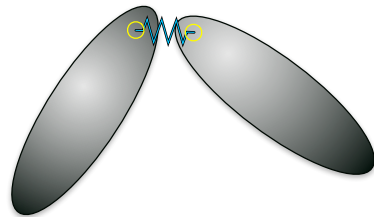
- A force / torque pair is a couple
  - Also a wrench
- Many couples are equivalent



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# Constraints

- Simplest method is to use spring attachments
  - Basically a penalty method



- Spring strength required to get good results may be unreasonably high
  - There are ways to cheat in some contexts...

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# Constraints

- Articulation constraints
  - Spring trick is an example of a full coordinate method
    - Better constraint methods exist
  - Reduced coordinate methods use DOFs in kinematic skeleton for simulation
    - Much more complex to explain
- Collisions
  - Penalty methods can also be used for collisions
  - Again, better constraint methods exist

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# Suggested Reading

- Brian Mirtich, "Fast and Accurate Computation of Polyhedral Mass Properties," Journal of Graphics Tools, volume 1, number 2, 1996. <http://www.cs.berkeley.edu/~jfc/mirtich/papers/vollnt.ps>
- Brian Mirtich and John Canny, "Impulse-based Simulation of Rigid Bodies," in Proceedings of 1995 Symposium on Interactive 3D Graphics, April 1995. <http://www.cs.berkeley.edu/~jfc/mirtich/papers/bsrb.ps>
- D. Baraff. Linear-time dynamics using Lagrange multipliers. Computer Graphics Proceedings, Annual Conference Series: 137-146, 1996. <http://www.pixar.com/companyinfo/research/deb/sig96.pdf>
- D. Baraff. Fast contact force computation for nonpenetrating rigid bodies. Computer Graphics Proceedings, Annual Conference Series: 23-34, 1994. <http://www.pixar.com/companyinfo/research/deb/sig94.pdf>

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Sunday, April 28, 13