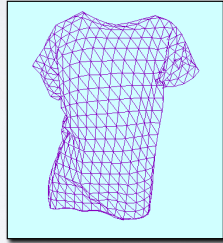
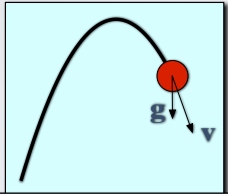


Physically Based Animation

3

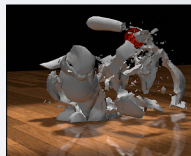
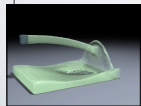
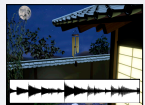
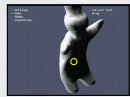
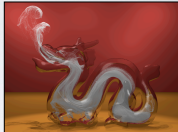
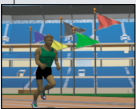
- Generate motion of objects using numerical simulation methods



$$\mathbf{x}^{t+\Delta t} = \mathbf{x}^t + \Delta t \mathbf{v}^t + \frac{1}{2} \Delta t^2 \mathbf{a}^t$$

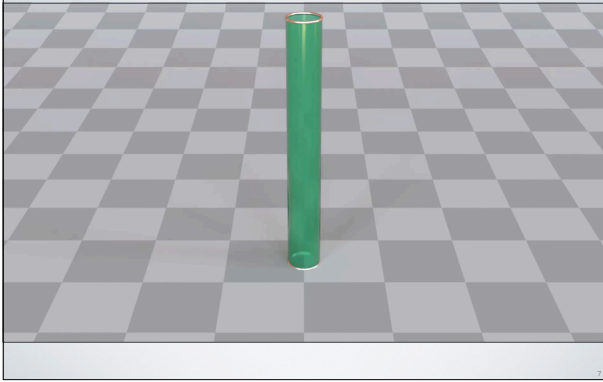
Physically Based Animation

4



Example: Fracture

7



Example: Fracture

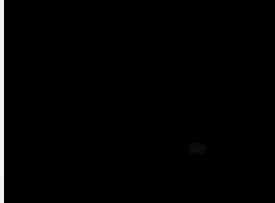
8



Particle Systems

11

- Single particles are very simple
- Large groups can produce interesting effects
- Supplement basic ballistic rules
 - Collisions
 - Interactions
 - Force fields
 - Springs
 - Others...



Feldman, Klingner, O'Brien, SIGGRAPH 2005

11

Basic Particles

12

- Basic governing equation
- \mathbf{f} is a sum of a number of things
 - Gravity: constant downward force proportional to mass
 - Simple drag: force proportional to negative velocity
 - Particle interactions: particles mutually attract and/or repel
 - Beware $O(n^2)$ complexity!
 - Force fields
 - Wind forces
 - User interaction

$$\ddot{\mathbf{x}} = \frac{1}{m} \mathbf{f}$$

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Basic Particles

13

- Properties other than position
 - Color
 - Temp
 - Age
- Differential equations also needed to govern these properties
- Collisions and other constrains directly modify position and/or velocity

13

Particle Rules

14



Multiple Burst

Bryan E. Feldman, James F. O'Brien, and Olan Arikun. "Animating Suspended Particle Explosions." In Proceedings of ACM SIGGRAPH 2003, pages 708-715, August 2003.

14

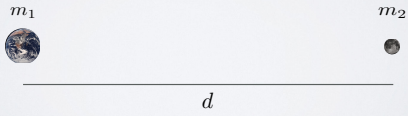
Gravitational Attraction

15

- Newton's universal law of gravitation
 - Gravitational pull between particles

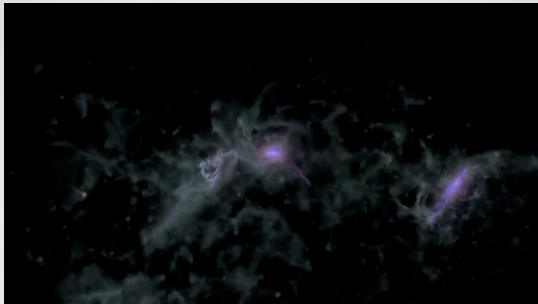
$$F_g = G \frac{m_1 m_2}{d^2}$$

$$G = 6.67428 \times 10^{-11} \text{ Nm}^2\text{kg}^{-2}$$



Example: Galaxy Simulation

16



Disk galaxy simulation, NASA Goddard

Integration

17

- Euler's Method
 - Simple
 - Commonly used
 - Very inaccurate
 - Most often goes unstable

$$\mathbf{x}^{t+\Delta t} = \mathbf{x}^t + \Delta t \dot{\mathbf{x}}^t$$

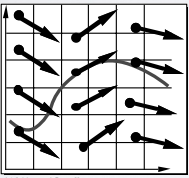
$$\dot{\mathbf{x}}^{t+\Delta t} = \dot{\mathbf{x}}^t + \Delta t \ddot{\mathbf{x}}^t$$

17

Integration

18

- For now let's pretend $\mathbf{f} = m\mathbf{v}$
 - *Velocity* (rather than acceleration) is a function of force



Wolken and Baraff

$$\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x}, t)$$

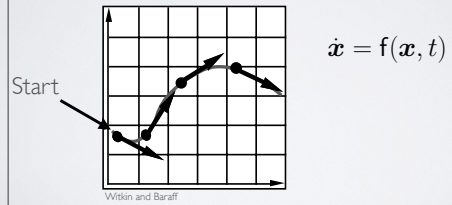
Note: Second order ODEs can be turned into first order ODEs using extra variables.

18

Integration

19

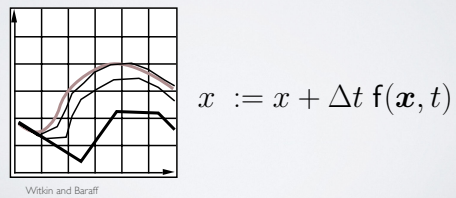
- For now let's pretend $f = mv$
- Velocity (rather than acceleration) is a function of force



Integration

20

- With numerical integration, errors accumulate
- Euler integration is particularly bad

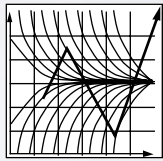
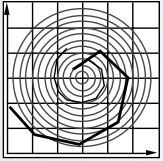


Integration

21

• Stability issues can also arise

- Occurs when errors lead to larger errors
- Often more serious than error issues



$$\dot{\mathbf{x}} = [-\sin(\omega t), -\cos(\omega t)]$$

Witkin and Baraff

21

Integration

22

• Modified Euler

$$\mathbf{x}^{t+\Delta t} = \mathbf{x}^t + \frac{\Delta t}{2} (\dot{\mathbf{x}}^t + \dot{\mathbf{x}}^{t+\Delta t})$$

$$\dot{\mathbf{x}}^{t+\Delta t} = \dot{\mathbf{x}}^t + \Delta t \ddot{\mathbf{x}}^t$$

$$\mathbf{x}^{t+\Delta t} = \mathbf{x}^t + \Delta t \dot{\mathbf{x}}^t + \frac{(\Delta t)^2}{2} \ddot{\mathbf{x}}^t$$

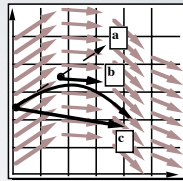
22

Integration

23

- Midpoint method

- Compute half Euler step
- Eval. derivative at halfway
- Retake step



Writkin and Baraff

23

Integration

24

- Implicit methods

- Informally (incorrectly) called backward methods
- Use derivatives in the future for the current step

$$\mathbf{x}^{t+\Delta t} = \mathbf{x}^t + \Delta t \dot{\mathbf{x}}^{t+\Delta t}$$

$$\dot{\mathbf{x}}^{t+\Delta t} = \dot{\mathbf{x}}^t + \Delta t \ddot{\mathbf{x}}^{t+\Delta t}$$

$$\dot{\mathbf{x}}^{t+\Delta t} = \mathbf{V}(\mathbf{x}^{t+\Delta t}, \dot{\mathbf{x}}^{t+\Delta t}, t + \Delta t)$$

$$\ddot{\mathbf{x}}^{t+\Delta t} = \mathbf{A}(\mathbf{x}^{t+\Delta t}, \dot{\mathbf{x}}^{t+\Delta t}, t + \Delta t)$$

24

Integration

25

• Implicit methods

- Informally (incorrectly) called backward methods
- Use derivatives in the future for the current step

$$\dot{\mathbf{x}}^{t+\Delta t} = \dot{\mathbf{x}}^t + \Delta t \mathbf{V}(\mathbf{x}^{t+\Delta t}, \dot{\mathbf{x}}^{t+\Delta t}, t + \Delta t)$$

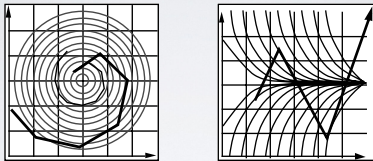
$$\dot{\mathbf{x}}^{t+\Delta t} = \dot{\mathbf{x}}^t + \Delta t \mathbf{A}(\mathbf{x}^{t+\Delta t}, \dot{\mathbf{x}}^{t+\Delta t}, t + \Delta t)$$

- Solve nonlinear problem for $\mathbf{x}^{t+\Delta t}$ and $\dot{\mathbf{x}}^{t+\Delta t}$
- This is fully implicit backward Euler
- Many other implicit methods exist...
- Modified Euler is *partially* implicit as is Verlet

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Temp Slide

26



Need to draw reverse diagrams...

26

Integration

27

- Semi-Implicit
 - Approximate with linearized equations

$$\mathbf{V}(\mathbf{x}^{t+\Delta t}, \dot{\mathbf{x}}^{t+\Delta t}) \approx \mathbf{V}(\mathbf{x}^t, \dot{\mathbf{x}}^t) + \mathbf{A} \cdot (\Delta \mathbf{x}) + \mathbf{B} \cdot (\Delta \dot{\mathbf{x}})$$

$$\mathbf{A}(\mathbf{x}^{t+\Delta t}, \dot{\mathbf{x}}^{t+\Delta t}) \approx \mathbf{A}(\mathbf{x}^t, \dot{\mathbf{x}}^t) + \mathbf{C} \cdot (\Delta \mathbf{x}) + \mathbf{D} \cdot (\Delta \dot{\mathbf{x}})$$

$$\begin{bmatrix} \mathbf{x}^{t+\Delta t} \\ \dot{\mathbf{x}}^{t+\Delta t} \end{bmatrix} = \begin{bmatrix} \mathbf{x}^t \\ \dot{\mathbf{x}}^t \end{bmatrix} + \Delta t \left(\begin{bmatrix} \dot{\mathbf{x}}^t \\ \ddot{\mathbf{x}}^t \end{bmatrix} + \begin{bmatrix} \mathbf{A} & \mathbf{B} \\ \mathbf{C} & \mathbf{D} \end{bmatrix} \begin{bmatrix} \Delta \mathbf{x} \\ \Delta \dot{\mathbf{x}} \end{bmatrix} \right)$$

27

Integration

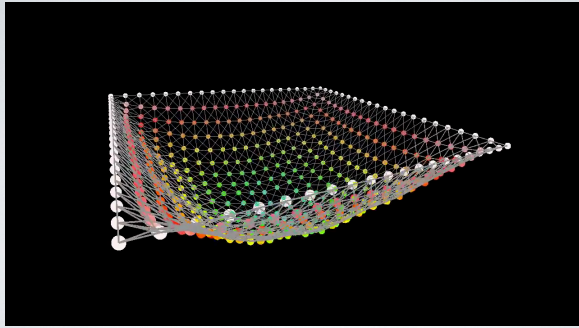
28

- Explicit methods can be conditionally stable
 - Depends on time-step and *stiffness* of system
- Fully implicit can be **un**conditionally stable
 - May still have large errors
- Semi-implicit can be conditionally stable
 - Nonlinearities can cause instability
 - Generally more stable than explicit
 - Comparable errors as explicit
 - Often show up as excessive damping

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Example: Mass Spring Mesh

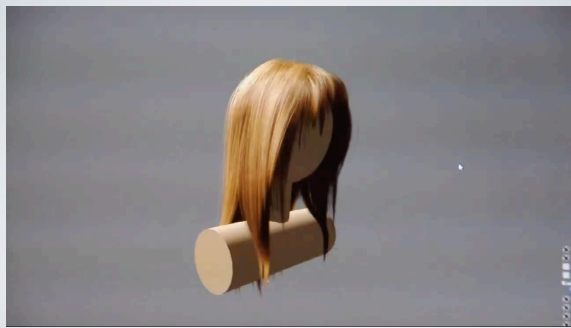
31



Slide from Ren Ng

Example: Hair

32



Slide from Ren Ng

Example: Cloth

33

Hanging Cloth

Huanmin Wang, Ravi Ramamoorthi, and James F. O'Brien. "Data-Driven Elastic Models for Cloth: Modeling and Measurement." ACM Transactions on Graphics, 30(5):7:1-11, July 2011. Proceedings of ACM SIGGRAPH 2011, Vancouver, BC Canada.

Example: Clothing on Character

34



Strain Limiting

35


Bunny
Hollow Triangle Mesh
59K Elements

Huanmin Wang, James F. O'Brien, and Ravi Ramamoorthi. "Multi-Resolution Isotropic Strain Limiting". In Proceedings of ACM SIGGRAPH Asia 2010, pages 160-1-10, December 2010.

A Simple Spring

36

- Ideal **zero**-length spring

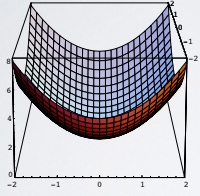
 $f_{a \rightarrow b} = k_s(\mathbf{b} - \mathbf{a})$

- Force pulls points together $f_{b \rightarrow a} = -f_{a \rightarrow b}$
- Strength proportional to distance

A Simple Spring

37

• Energy potential



$$E = 1/2 k_S (\mathbf{b} - \mathbf{a}) \cdot (\mathbf{b} - \mathbf{a})$$

$$\mathbf{f}_{a \rightarrow b} = k_S (\mathbf{b} - \mathbf{a})$$

$$\mathbf{f}_{b \rightarrow a} = -\mathbf{f}_{a \rightarrow b}$$

$$\mathbf{f}_a = -\nabla_a E = - \left[\frac{\partial E}{\partial a_x}, \frac{\partial E}{\partial a_y}, \frac{\partial E}{\partial a_z} \right]$$

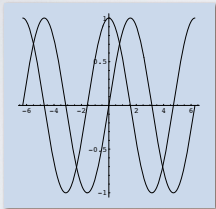


37

A Simple Spring

38

• Energy potential: kinetic vs elastic



$$E = 1/2 k_S (\mathbf{b} - \mathbf{a}) \cdot (\mathbf{b} - \mathbf{a})$$

$$E = 1/2 m (\dot{\mathbf{b}} - \dot{\mathbf{a}}) \cdot (\dot{\mathbf{b}} - \dot{\mathbf{a}})$$



38

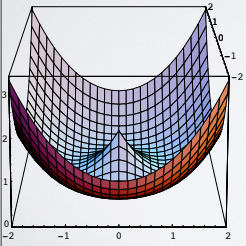
Non-Zero Length Springs

39



$$\mathbf{f}_{a \rightarrow b} = k_s \frac{\mathbf{b} - \mathbf{a}}{\|\mathbf{b} - \mathbf{a}\|} (\|\mathbf{b} - \mathbf{a}\| - l)$$

Rest length



$$E = k_s (\|\mathbf{b} - \mathbf{a}\| - l)^2$$

Comments on Springs

40


- Springs with zero rest length are linear
- Springs with non-zero rest length are nonlinear
 - Force **magnitude** linear w/ displacement (from rest length)
 - Force direction is non-linear
 - Singularity at

$$\|\mathbf{b} - \mathbf{a}\| = 0$$

Damping

41

- "Mass proportional" damping


$$\mathbf{f} = -k_d \dot{\mathbf{a}}$$


- Behaves like viscous drag on all motion
- Consider a pair of masses connected by a spring
 - How to model rusty vs oiled spring
 - Should internal damping slow group motion of the pair?
- Can help stability... up to a point

41

Damping

42

- "Stiffness proportional" damping


$$\mathbf{f}_a = -k_d \frac{\mathbf{b} - \mathbf{a}}{\|\mathbf{b} - \mathbf{a}\|^2} (\mathbf{b} - \mathbf{a}) \cdot (\dot{\mathbf{b}} - \dot{\mathbf{a}})$$

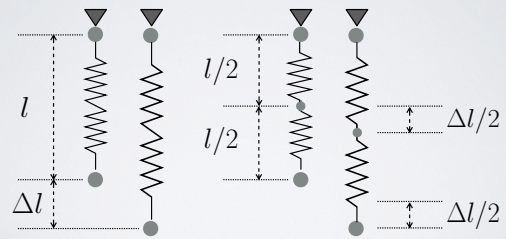
- Behaves viscous drag on change in spring length
- Consider a pair of masses connected by a spring
 - How to model rusty vs oiled spring
 - Should internal damping slow group motion of the pair?

42

Spring Constants

43

- Two ways to model a single spring



Spring Constants

44

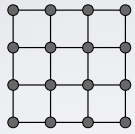
- Constant k_s gives inconsistent results with different discretizations
- Change in length is not what we want to measure
- Strain: change in length as fraction of original length

$$\epsilon = \frac{\Delta l}{l_0} \quad \text{Nice and simple for 1D...}$$

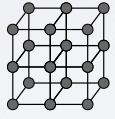
Structures from Springs

45

- Sheets



- Blocks



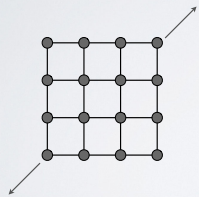
- Others

45

Structures from Springs

46

- They behave like what they are (obviously!)



This structure will not resist shearing

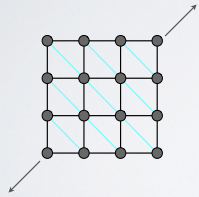
This structure will not resist out-of-plane bending either...

46

Structures from Springs

47

- They behave like what they are (obviously!)



This structure will resist shearing
but has anisotropic bias

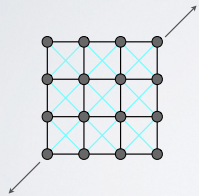
This structure still will not resist
out-of-plane bending

47

Structures from Springs

48

- They behave like what they are (obviously!)



This structure will resist shearing
Less bias

Interference between spring sets

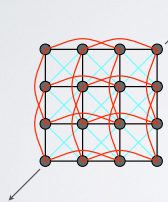
This structure still will not resist
out-of-plane bending

48

Structures from Springs

49

- They behave like what they are (obviously!)



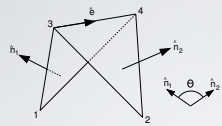
This structure will resist shearing
Less bias
Interference between spring sets

This structure will resist out-of-plane bending
Interference between spring sets
Odd behavior

How do we set spring constants? ..

Edge Springs

50



$$u_1 = |E| \frac{N_1}{|N_1|^2} \quad u_2 = |E| \frac{N_2}{|N_2|^2}$$

$$u_3 = \frac{(x_1 - x_4) \cdot E}{|E|} \frac{N_1}{|N_1|^2} + \frac{(x_2 - x_4) \cdot E}{|E|} \frac{N_2}{|N_2|^2}$$

$$u_4 = -\frac{(x_1 - x_3) \cdot E}{|E|} \frac{N_1}{|N_1|^2} - \frac{(x_2 - x_3) \cdot E}{|E|} \frac{N_2}{|N_2|^2}$$

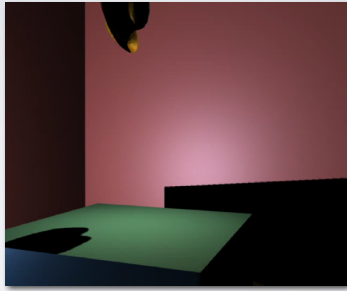
$$F_i^e = k^e \frac{|E|^2}{|N_1| + |N_2|} \sin(\theta/2) u_i$$

From Bridson *et al.*, 2003, also see Grinspun *et al.*, 2003

50

Example: Thin Material

51



Discrete Shells
SCA 2003
Eitan Grinspun, Anil Hirani, Mathieu Desbrun and Peter Schröder

Sharp Creases

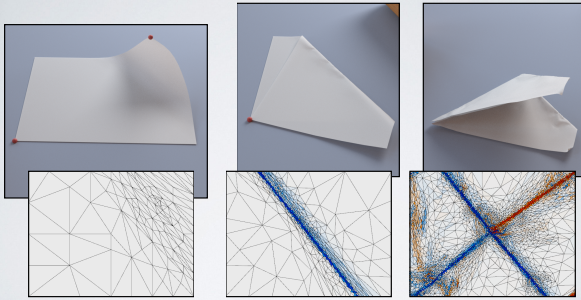
52



Paper Folding

Sharp Creases

53



Anisotropic remeshing avoids locking.

Fracture

54



Dominated by discretization artifacts



Natural appearance

