

CS-184: Computer Graphics

Lecture #23: Rigid Body Dynamics

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V2011F-23-1.0

Announcement

- Final Project Poster Session
 - Thursday, Friday December 8th, 2:30-5:00 pm
 - Poster stands and tables provided
 - Laptop videos or demos are highly recommended
 - Limited AC outlets
- Final project reports
 - **Hardcopy** due to me by December 15th 5pm.
 - *No time for late submission!*
- Final exam
 - Tuesday, December 13th, 8:00 - 11:00 am
 - 10 Evans

Today

- Rigid-body dynamics
- Articulated systems

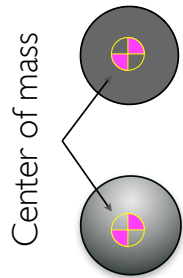
3

A Rigid Body

- **A solid object that does not deform**
- Consists of infinite number of infinitesimal mass points...
- ...that share a single RB transformation
 - Rotation + Translation (no shear or scale)
- $$\mathbf{x}^W = \mathbf{R} \cdot \mathbf{x}^L + \mathbf{t}$$
- Rotation and translation vary over time
- Limit of deformable object as $k_s \rightarrow \infty$

3

A Rigid Body



In 2D:
Translation 2 “directions”
Rotation 1 “direction”
3 DOF Total

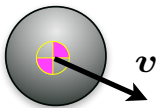
In 3D:
Translation 3 “directions”
Rotation 3 “direction”
6 DOF Total

Translation and rotation are **decoupled**

2D is boring... we'll stick to 3D from now on...

5

Translational Motion



Just like a point mass:

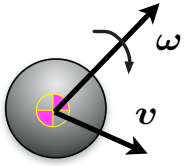
$$\dot{\mathbf{p}} = \mathbf{v}$$

$$\dot{\mathbf{v}} = \mathbf{a} = \mathbf{f}/m$$

Note: Recall discussion on integration...

6

Rotational Motion



Rotation gets a bit odd, as well see...

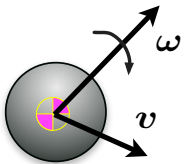
Rotational "position" R

Rotation matrix
Exponential map
Quaternions

Rotational velocity ω
Stored as a vector
(Also called angular velocity...)
Measured in radians / second

7

Rotational Motion



Kinetic energy due to rotation:

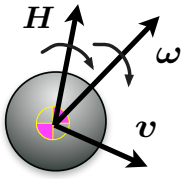
"Sum energy (from rotation) over all points in the object"

$$E = \int_{\Omega} \frac{1}{2} \rho \dot{\mathbf{x}} \cdot \dot{\mathbf{x}} du$$

$$E = \int_{\Omega} \frac{1}{2} \rho ([\omega \times] \mathbf{x}) \cdot ([\omega \times] \mathbf{x}) du$$

8

Rotational Motion



Angular momentum
 Similar to linear momentum
 Can be derived from rotational energy

Figure is a lie if this really is a sphere...

$$\mathbf{H} = \int_{\Omega} \rho \mathbf{x} \times \dot{\mathbf{x}} \, du$$

$$\mathbf{H} = \int_{\Omega} \rho \mathbf{x} \times (\boldsymbol{\omega} \times \mathbf{x}) \, du$$

$$\mathbf{H} = \left(\int_{\Omega} \dots \, du \right) \boldsymbol{\omega}$$

"Inertia Tensor" not identity matrix...

$$\mathbf{H} = \mathbf{I} \boldsymbol{\omega}$$

Rotational Motion

$H = \frac{\partial E}{\partial \omega}$ *H momentum (angular) (work/rot)*

$H_p = \frac{\partial E}{\partial \omega_p}$

$$= \int_{\Omega} \rho \mathbf{x} \left(\epsilon_{ijk} \delta_{jp} x_k \epsilon_{iab} \omega_a x_b + \epsilon_{ijn} \omega_j x_n \epsilon_{iab} \delta_{pa} x_b \right) du$$

$$= \int_{\Omega} \rho \mathbf{x} \left(\epsilon_{ijk} x_n \epsilon_{iab} \omega_a x_b + \epsilon_{ijn} \omega_j x_n \epsilon_{iab} x_b \right) du$$

$$= \int_{\Omega} \rho \epsilon_{ijk} x_n \epsilon_{iab} \omega_a x_b \, du$$

$$= \int_{\Omega} \rho \mathbf{x} \times (\boldsymbol{\omega} \times \mathbf{x}) \, du$$

$$= \int_{\Omega} \rho \mathbf{x} \times \mathbf{v} \, du$$

$$= \omega_a \int_{\Omega} \rho \epsilon_{ijk} x_n \epsilon_{iab} x_b \, du$$

$$= \omega_a \int_{\Omega} \rho (\delta_{pa} \delta_{kn} - \delta_{pn} \delta_{ka}) x_n x_b \, du$$

$$= \omega_a \int_{\Omega} \rho (\delta_{pa} x_n x_n - x_n x_p) \, du$$

* $\int_{\Omega} H_p = \int_{\Omega} \rho \omega_a \, du$ *Inertia Tensor*

Angular momentum
 Similar to linear momentum
 Can be derived from rotational energy

$$\mathbf{H} = \int_{\Omega} \rho \mathbf{x} \times \dot{\mathbf{x}} \, du$$

$$\mathbf{H} = \int_{\Omega} \rho \mathbf{x} \times (\boldsymbol{\omega} \times \mathbf{x}) \, du$$

$$\mathbf{H} = \left(\int_{\Omega} \dots \, du \right) \boldsymbol{\omega}$$

"Inertia Tensor" not identity matrix...

$$\mathbf{H} = \mathbf{I} \boldsymbol{\omega}$$

Rotational Motion

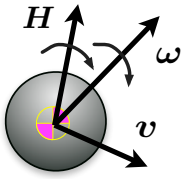


Figure is a lie if this really is a sphere...

Angular momentum
Similar to linear momentum
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$$\mathbf{H} = \int_{\Omega} \rho \mathbf{x} \times \dot{\mathbf{x}} du$$

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"Inertia Tensor" not identity matrix...

$$\mathbf{H} = \mathbf{I}\boldsymbol{\omega}$$

9

Inertia Tensor

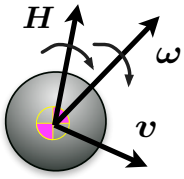
$$\mathbf{I} = \int_{\Omega} \rho \begin{bmatrix} y^2 + z^2 & -xy & -xz \\ -xy & z^2 + x^2 & -yz \\ -xz & -yz & x^2 + y^2 \end{bmatrix} du$$

See example for simple shapes at
<http://scienceworld.wolfram.com/physics/MomentofInertia.html>

Can also be computed from polygon models by transforming volume integral to a surface one.
See paper/code by Brian Mirtich.

10

Rotational Motion



Conservation of momentum:

$$\mathbf{H}^W = \mathbf{I}^W \boldsymbol{\omega}^W$$

$$\mathbf{H}^W = \mathbf{R} \mathbf{I}^L \mathbf{R}^T \boldsymbol{\omega}^W$$

$$\dot{\mathbf{H}}^W = \dot{\mathbf{R}} \mathbf{I}^L \mathbf{R}^T \boldsymbol{\omega}^W + \cancel{\mathbf{R} \mathbf{I}^L \dot{\mathbf{R}}^T \boldsymbol{\omega}^W} + \mathbf{R} \mathbf{I}^L \mathbf{R}^T \boldsymbol{\alpha}^W$$

Figure is a lie if this really is a sphere...

$$\dot{\mathbf{H}}^W = 0$$

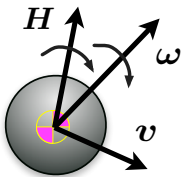
$$\dot{\mathbf{R}} = \boldsymbol{\omega} \times \mathbf{R}$$

$$\boldsymbol{\alpha}^W = (\mathbf{R} \mathbf{I}^L \mathbf{R}^T)^{-1} (-\boldsymbol{\omega}^W \times \mathbf{H}^W)$$

In other words, things wobble when they rotate.

11

Rotational Motion



$$\dot{\mathbf{R}} = [\boldsymbol{\omega} \times] \mathbf{R}$$

$$\dot{\boldsymbol{\omega}} = \boldsymbol{\alpha}$$

Figure is a lie if this really is a sphere...

$$\boldsymbol{\alpha}^W = (\mathbf{R} \mathbf{I}^L \mathbf{R}^T)^{-1} ((-\boldsymbol{\omega}^W \times \mathbf{H}^W) + \boldsymbol{\tau})$$

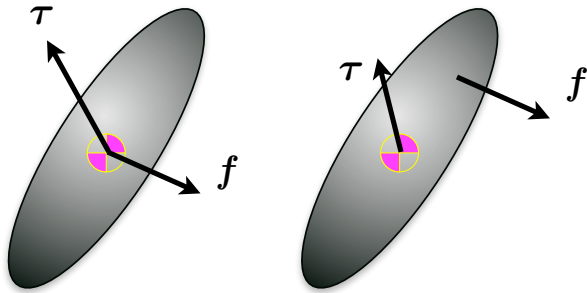
$$\boldsymbol{\tau} = \mathbf{f} \times \mathbf{x}$$

Take care when integrating rotations, they need to stay rotations.

12

Couples

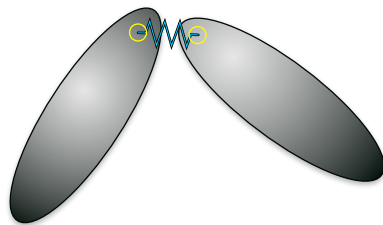
- A force / torque pair is a couple
 - Also a wrench
- Many couples are equivalent



13

Constraints

- Simplest method is to use spring attachments
 - Basically a penalty method



- Spring strength required to get good results may be unreasonably high
 - There are ways to cheat in some contexts...

14

Constraints

- Articulation constraints
 - Spring trick is an example of a full coordinate method
 - Better constraint methods exist
 - Reduced coordinate methods use DOFs in kinematic skeleton for simulation
 - Much more complex to explain
- Collisions
 - Penalty methods can also be used for collisions
 - Again, better constraint methods exist

Suggested Reading

• Brian Mirtich, "Fast and Accurate Computation of Polyhedral Mass Properties," Journal of Graphics Tools, volume 1, number 2, 1996. <http://www.cs.berkeley.edu/~jfc/mirtich/papers/vollnt.ps>

• Brian Mirtich and John Canny, "Impulse-based Simulation of Rigid Bodies," in Proceedings of 1995 Symposium on Interactive 3D Graphics, April 1995. <http://www.cs.berkeley.edu/~jfc/mirtich/papers/ibsr.ps>

• D. Baraff, Linear-time dynamics using Lagrange multipliers. Computer Graphics Proceedings, Annual Conference Series: 137-146, 1996. <http://www.pixar.com/companyinfo/research/deb/sig96.pdf>

• D. Baraff, Fast contact force computation for nonpenetrating rigid bodies. Computer Graphics Proceedings, Annual Conference Series: 23-34, 1994. <http://www.pixar.com/companyinfo/research/deb/sig94.pdf>